

Kinetic Theory of Gases

(NEP Semester IV - Chapter 3)

Dr. DEBOJYOTI HALDER

*Department of Physics,
R. B. C. Evening College*

Ideal Gas Equation

Boyle's Law

states :

“The pressure of a fixed mass of gas at constant temperature is inversely proportional to its volume.”

OR

$$P \propto \frac{1}{V} \quad \text{if } T = \text{constant}$$

- The related equations to the Boyle's law are

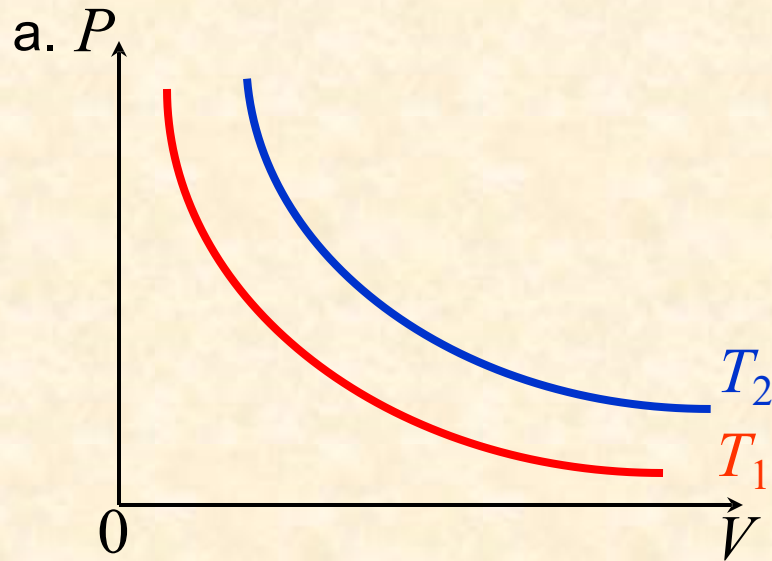
$$PV = \text{constant}$$

OR

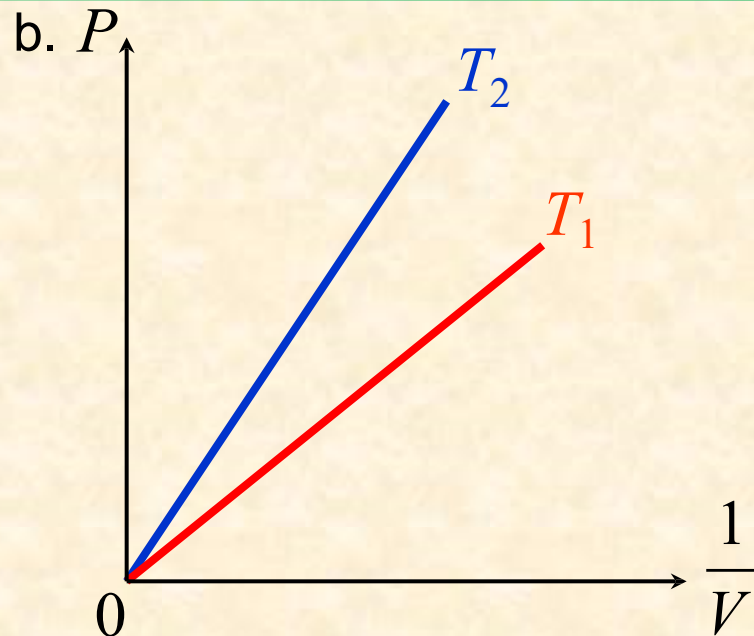
$$P_1V_1 = P_2V_2$$

where P_1 : initial pressure
 P_2 : final pressure
 V_1 : initial volume
 V_2 : final volume

- Graphs of the Boyle's law.



The pressure of a fixed mass of gas at constant temperature is inversely proportional to its volume.



Charles's law

- states:

“The volume of a fixed mass of gas at constant pressure is directly proportional to its absolute temperature.”

OR

$$V \propto T \quad \text{if } P = \text{constant}$$

The related equations to the Charles' law are

$$\frac{V}{T} = \text{constant}$$

OR

$$\frac{V_1}{T_1} = \frac{V_2}{T_2}$$

where

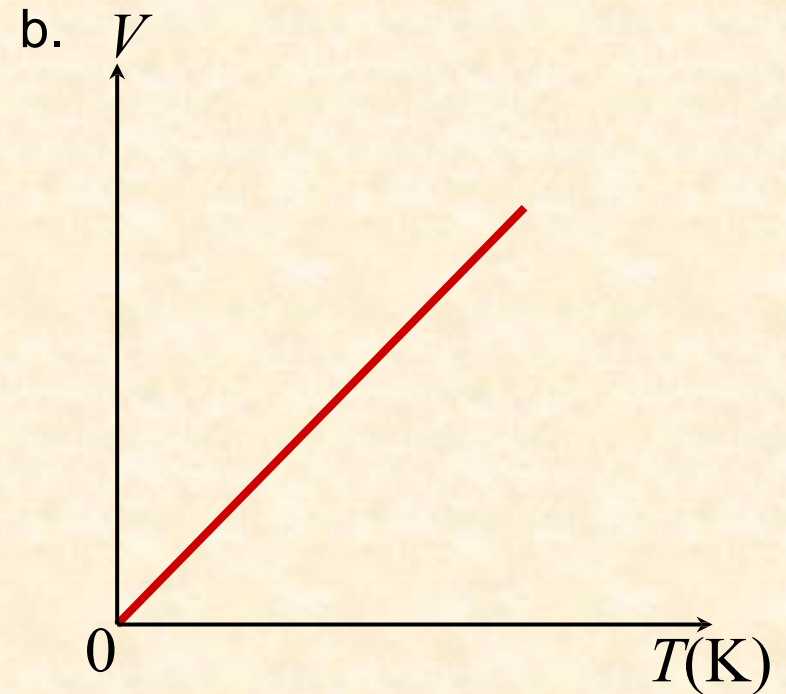
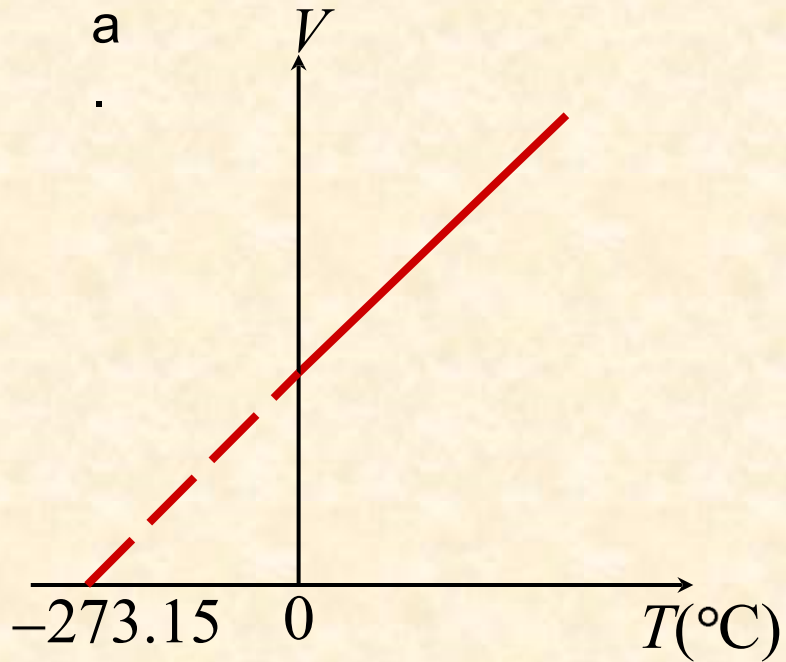
T_1 : initial absolute temperature

T_2 : final absolute temperature

V_1 : initial volume

V_2 : final volume

Graphs of the Charles's law.



Gay-Lussac's (pressure) law
states :

“The pressure of a fixed mass of gas at constant volume is directly proportional to its absolute temperature.”

OR

$$P \propto T \text{ if } V = \text{constant}$$

The related equations to the Gay-Lussac's law are

$$\frac{P}{T} = \text{constant}$$

OR

$$\frac{P_1}{T_1} = \frac{P_2}{T_2}$$

where

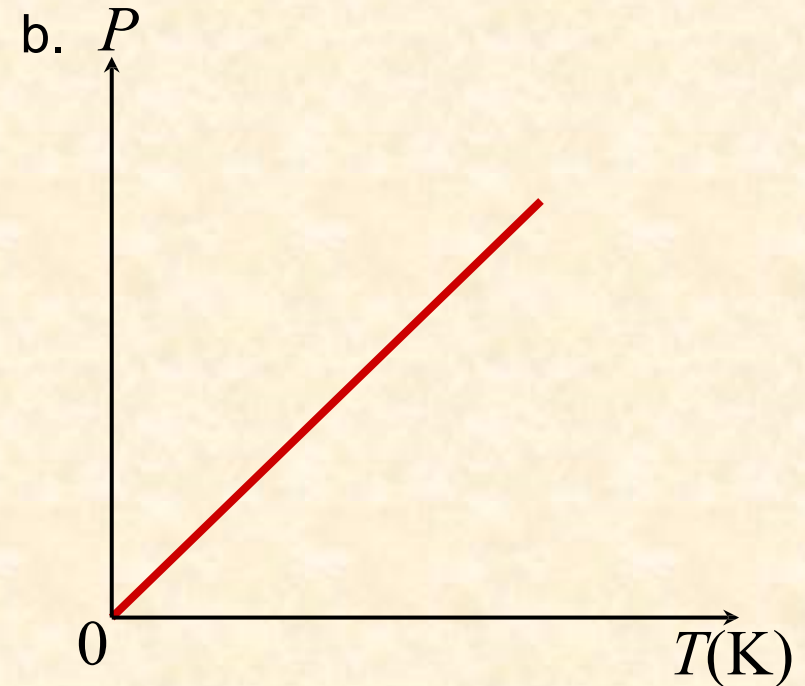
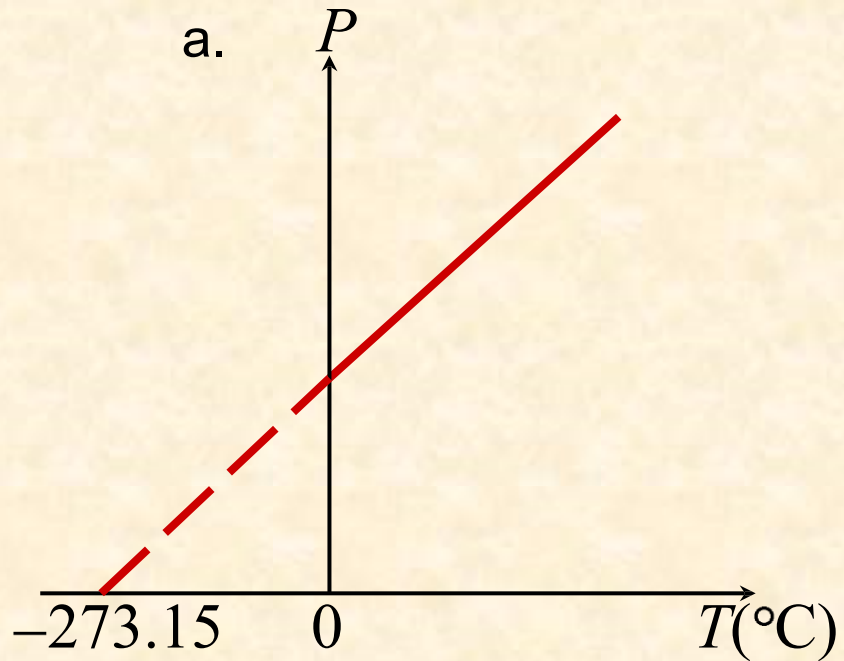
T_1 : initial absolute temperature

T_2 : final absolute temperature

P_1 : initial pressure

P_2 : final pressure

Graphs of the Gay-lussac's (pressure) law



Equation of state for an ideal gas

- An **ideal gas** is defined as **a perfect gas which obeys the three gas laws (Boyle's, Charles's and Gay-Lussac's) exactly.**
- Consider an ideal gas in a container changes its pressure P , volume V and temperature T as shown in Figure 1.

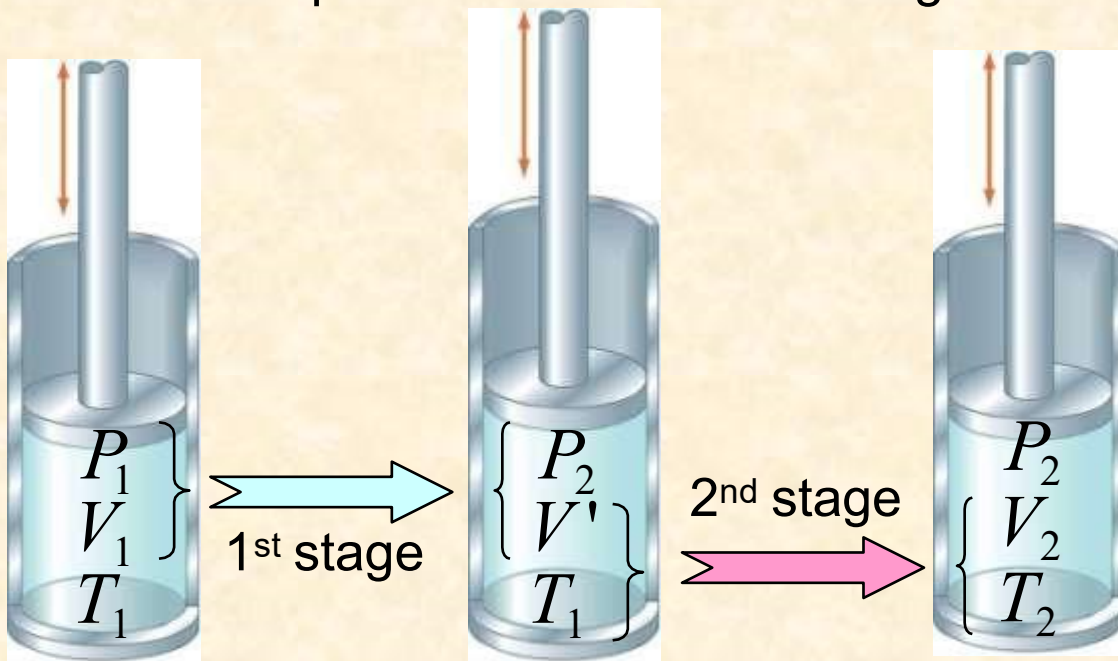


Figure 1

- In 1st stage, temperature is kept at T_1 ,
Using Boyle's law :

$$P_2 V' = P_1 V_1$$

$$\Rightarrow V' = \frac{P_1 V_1}{P_2} \quad (1)$$

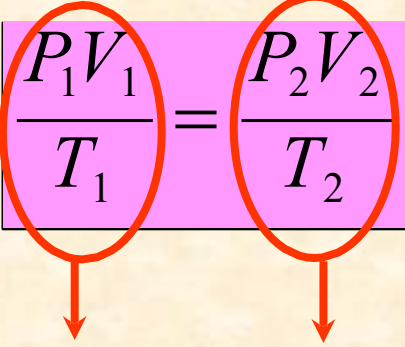
- In 2nd stage, pressure is kept constant at P_2 ,
Using Charles's law :

$$\frac{V'}{T_1} = \frac{V_2}{T_2}$$

$$\Rightarrow V' = \frac{V_2 T_1}{T_2} \quad (2)$$

- Equating eqs. (1) and (2), thus

$$\frac{P_1 V_1}{T_1} = \frac{P_2 V_2}{T_2} \quad \text{OR}$$


Initial Final

$$\frac{PV}{T} = \text{constant} \quad (3)$$

- Consider **1 mole** of gas at **standard temperature and pressure (S.T.P.)**, $T = 273.15 \text{ K}$, $P = 101.3 \text{ kPa}$ and $V_m = 0.0224 \text{ m}^3$

- From equation (3),

$$R = \frac{PV_m}{T} = \frac{101.3 \times 10^3 \times 0.0224}{273.15}$$
$$R = 8.31 \text{ J K}^{-1} \text{ mol}^{-1}$$

where **R** is called **molar gas constant** and its **value** is the **same for all gases**.

- Thus $\frac{PV_m}{T} = R$

$$PV_m = RT \text{ where } V_m : \text{volume of 1 mole gas}$$

- For **n mole** of an **ideal gas**, the **equation of state** is written as

$$PV = nRT$$

- where n : the number of mole gas

$$n = \frac{m}{M}$$

where

m : mass of a gas

M : molar mass of a gas

OR

$$n = \frac{N}{N_A}$$

where

N : number of molecules

N_A : Avogadro's constant

$$= 6.02 \times 10^{23} \text{ mol}^{-1}$$

– If the Boltzmann constant, k is defined as

$$k = \frac{R}{N_A} = 1.38 \times 10^{-23} \text{ J K}^{-1}$$

then the equation of state becomes

$$PV = NkT$$

Example 1 :

The volume of vessel A is three times of the volume vessel B. The vessels are filled with an ideal gas and are at a steady state. The temperature of vessel A and vessel B are at 300 K and 500 K respectively as shown in Figure 2.

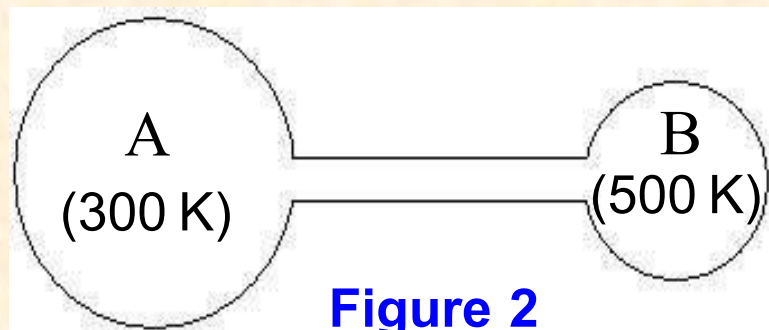


Figure 2

If the mass of the gas in the vessel A is m , obtain the mass of the gas in the vessel B in terms of m .

Solution :

$$V_A = 3V_B; m_A = m; T_{0A} = 300 \text{ K}; T_{0B} = 500 \text{ K}$$

Since the vessels A and B are connected by a narrow tube thus the pressure for both vessels is same, finally i.e.

$$P_A = P_B = P$$

The system is in the steady state, thus

$$T_{0A} = T_A = 300 \text{ K}; T_{0B} = T_B = 500 \text{ K}$$

By applying the equation of state for an ideal gas,

$$PV = nRT \text{ and } n = \frac{m}{M}$$

$$PV = \left(\frac{m}{M} \right) RT$$

Therefore,
Vessel A :

$$P_A V_A = \left(\frac{m_A}{M} \right) R T_A$$

$$P V_B = \left(\frac{m}{M} \right) R 100$$

Vessel B :

$$V_A = 3V_B; m_A = m; T_{0A} = 300 \text{ K}; T_{0B} = 500 \text{ K}$$

$$P V_B = 100 \left(\frac{m}{M} \right) R \quad \text{————— (1)}$$

$$P V_B = 500 \left(\frac{m_B}{M} \right) R \quad \text{————— (2)}$$

By equating the eqs. (1) and (2) hence

$$100 \left(\frac{m}{M} \right) R = 500 \left(\frac{m_B}{M} \right) R$$

$$m_B = \frac{m}{5}$$

2. Kinetic theory of gases

- The macroscopic behaviour of an ideal gas can be describe by using the equation of state but the **microscopic behaviour** only can be describe by **kinetic theory of gases**.

1 Assumption of kinetic theory of gases

- All gases are made up of **identical atoms or molecules**.
- All atoms or molecules move **randomly and haphazardly**.
- The **volume** of the atoms or molecules is **negligible** when **compared** with the **volume occupied** by the gas.
- The **intermolecular forces** are **negligible** except during **collisions**.
- Inter-atomic or molecular **collisions** are **elastic**.
- The **duration of a collision** is **negligible** compared with the **time spent travelling** between collisions.
- Atoms and molecules move with **constant speed** between **collisions**. **Gravity** has **no effect** on molecular motion.

2 Force exerted by an ideal gas

- Consider an ideal gas of N molecules are contained in a cubical container of side d as shown in Figure 4.

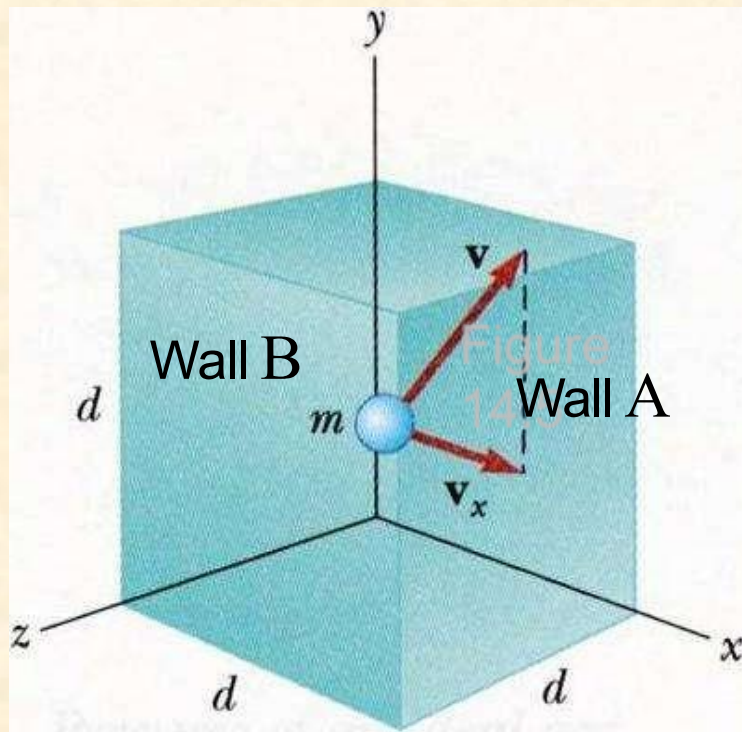


Figure 4

- Let each molecule of the gas have the mass m and velocity v .
- The velocity, v of each molecule can be resolved into their components i.e. v_x , v_y and v_z .
- Consider, initially a single molecule moving with a velocity v_x towards wall A and after colliding elastically, it moves in the opposite direction with a velocity $-v_x$ as shown in Figure 5.

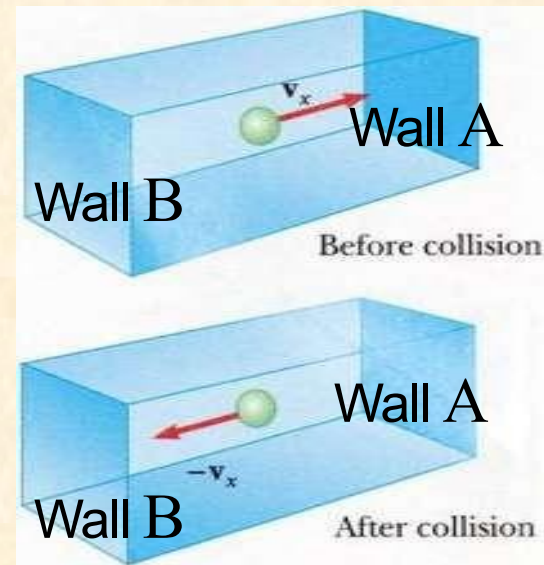


Figure 5

Therefore the change in the linear momentum of the molecule is given by

$$\Delta P_x = -mv_x - mv_x$$

$$\Delta P_x = -2mv_x$$

The molecule has to travel a distance $2d$ (from A to B and back to A) before its next collision with wall A. The time taken for this movement is

$$\Delta t = \frac{2d}{v_x}$$

If F_{x1} is the magnitude of the average force exerted by a molecule on the wall in the time Δt , thus by applying Newton's second law of motion gives

$$F_{x1} = \frac{\Delta P_x}{\Delta t} = \frac{2mv_x}{\left(\frac{2d}{v_x}\right)} \Rightarrow F_{x1} = \left(\frac{m}{d}\right)v_x^2$$

For N molecules of the ideal gas,

$$F_x = \left(\frac{m}{d}\right)v_{x1}^2 + \left(\frac{m}{d}\right)v_{x2}^2 + \dots + \left(\frac{m}{d}\right)v_{xN}^2$$

$$F_x = \frac{m}{d}v_{x1}^2 + v_{x2}^2 + \dots + v_{xN}^2$$

where v_{x1} is the x component of velocity of molecule 1, v_{x2} is the component of velocity molecule 2 and so on.

- The mean (average) value of the square of the velocity in the x direction for N molecules is

$$\langle v_x^2 \rangle = \frac{v_{x1}^2 + v_{x2}^2 + \dots + v_{xN}^2}{N}$$

$$v_{x1}^2 + v_{x2}^2 + \dots + v_{xN}^2 = N \langle v_x^2 \rangle$$

Thus, the x component for the total force exerted on the wall of the cubical container is

$$F_x = \frac{m}{d} N \langle v_x^2 \rangle$$

- The magnitude of the velocity v is given by

$$v^2 = v_x^2 + v_y^2 + v_z^2$$

then

$$\langle v^2 \rangle = \langle v_x^2 \rangle + \langle v_y^2 \rangle + \langle v_z^2 \rangle$$

- Since the velocities of the molecules in the ideal gas are **completely random**, there is **no preference to one direction or another**. Hence

$$\langle v_x^2 \rangle = \langle v_y^2 \rangle = \langle v_z^2 \rangle$$

$$\langle v^2 \rangle = 3 \langle v_x^2 \rangle$$

$$\langle v_x^2 \rangle = \frac{\langle v^2 \rangle}{3}$$

- The total force exerted on the wall in all direction, F is given by

$$F = \frac{m}{d} N \langle v_x^2 \rangle$$



$$F = \frac{m}{d} N \left(\frac{\langle v^2 \rangle}{3} \right)$$

$$F = \frac{N}{3} \left(\frac{m \langle v^2 \rangle}{d} \right)$$

where $\langle v^2 \rangle$: mean square velocity of the molecule

3 Pressure of an ideal gas

- From the definition of pressure,

$$P = \frac{F}{A} \text{ where } A = d^2 \text{ and } F = \frac{N}{3} \left(\frac{m \langle v^2 \rangle}{d} \right)$$

$$P = \frac{1}{3} \left(\frac{Nm \langle v^2 \rangle}{d^3} \right) \text{ and } d^3 = V$$

$$P = \frac{1}{3} \left(\frac{Nm}{V} \langle v^2 \rangle \right) \text{-----(1)}$$

$$PV = \frac{1}{3} Nm \langle v^2 \rangle \text{-----(2)}$$

where Nm : mass of an ideal gas in the container

- Since the density of the gas, ρ is given by

$$\rho = \frac{Nm}{V}$$

hence the equation (15.1) can be written as

$$P = \frac{1}{3} \rho \langle v^2 \rangle \quad \text{--- (3)}$$

where

P : pressure by the gas

ρ : density of the gas

$\langle v^2 \rangle$: mean square velocity of the gas molecules

4 Root mean square velocity (v_{rms})

- is defined as

$$v_{\text{rms}} = \sqrt{\langle v^2 \rangle}$$

- From the equation of state in terms of Boltzmann constant, k :

$$PV = NkT \quad \text{--- (4)}$$

- By equating the eqs. (15.4) and (14.2), thus

$$NkT = \frac{1}{3} Nm \langle v^2 \rangle$$

$$\langle v^2 \rangle = \frac{3kT}{m}$$

$$v_{\text{rms}} = \sqrt{\frac{3kT}{m}} \quad \text{OR} \quad v_{\text{rms}} = \sqrt{\frac{3RT}{M}}$$

where

v_{rms} : root meansquare velocity (speed)

m : mass of a molecule gas

M : molar mass of a gas

T : absolute temperature

- Since

$$P = \frac{1}{3} \rho \langle v^2 \rangle$$

thus

$$\langle v^2 \rangle = \frac{3P}{\rho}$$

therefore the equation of root mean square velocity of the gas molecules also can be written as

$$v_{\text{rms}} = \sqrt{\frac{3P}{\rho}}$$

Molecular kinetic energy and internal energy

1. Translational kinetic energy of molecule

- From equation (1), thus

$$P = \frac{1}{3} \left(\frac{Nm}{V} \langle v^2 \rangle \right)$$

$$P = \frac{2}{3} \left(\frac{N}{V} \right) \left(\frac{1}{2} m \langle v^2 \rangle \right) \dots\dots\dots (5)$$

This equation shows that

P increases (\uparrow) When



$\left(\frac{N}{V} \right)$ increases



$\left(\frac{1}{2} m \langle v^2 \rangle \right)$ increases

1 Translational kinetic energy of molecule

This equation shows that

P increases (\uparrow) When



$\left(\frac{N}{V}\right)$ increases



$\left(\frac{1}{2}m < v^2 >\right)$ increases

Rearrange equation (14.5), thus

$$PV = \frac{2}{3} N \left(\frac{1}{2} m < v^2 > \right)$$

and

$$PV = NkT$$

$$NkT = \frac{2}{3} N \left(\frac{1}{2} m \langle v^2 \rangle \right)$$

$$\frac{1}{2} m \langle v^2 \rangle = \frac{3}{2} kT$$

and

$$\frac{1}{2} m \langle v^2 \rangle = K_{\text{tr}}$$

$$K_{\text{tr}} = \frac{3}{2} kT = \frac{3}{2} \left(\frac{R}{N_A} \right) T$$

where

K_{tr} : average translational kinetic energy of a molecule

T : absolute temperature

k : Boltzmann constant

R : molar gas constant

N_A : Avogadro constant

- For N molecules of an ideal gas in the cubical container, the total average (mean) translational kinetic energy, E is given by

$$E = NK_{tr}$$

$$E = N \left(\frac{3}{2} kT \right)$$

$$E = \frac{3}{2} NkT$$

OR

$$E = \frac{3}{2} nRT$$

2 Principle of equipartition of energy

- States : “the mean (average) kinetic energy of every degrees of freedom of a molecule is $\frac{1}{2}kT$.”

Therefore

$$\langle K \rangle = \frac{f}{2} kT$$

➡ Mean (average) kinetic energy per molecule

OR

$$\langle K \rangle = \frac{f}{2} RT$$

➡ Mean (average) kinetic energy per mole

where f : degrees of freedom

T : absolute temperature

3 Degree of freedom (f)

is defined as a **number of independent ways in which an atom or molecule can absorb or release or store the energy.**

Monatomic gas (e.g. He, Ne, Ar)

- The number of degrees of freedom is **3** i.e. **three direction of translational motion** where contribute translational kinetic energy as shown in Figure 6.

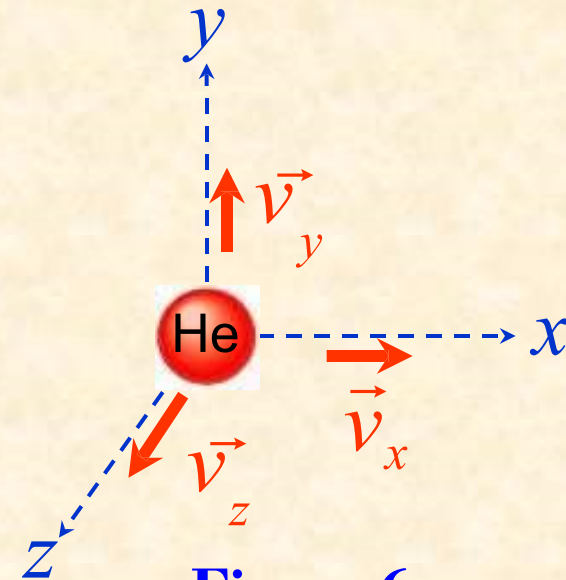


Figure 6

Diatomic gas (e.g. H_2 , O_2 , N_2)

- The number of degrees of freedom is

Translational kinetic energy $\rightarrow 3$

Rotational kinetic energy $\rightarrow \underline{2}$
5

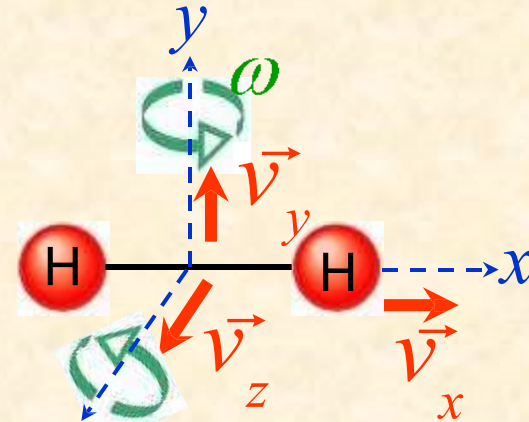


Figure 7

Polyatomic gas (e.g. H_2O , CO_2 , NH_3)

- The number of degrees of freedom is

Translational kinetic energy $\rightarrow 3$

Rotational kinetic energy $\rightarrow \underline{3}$
6

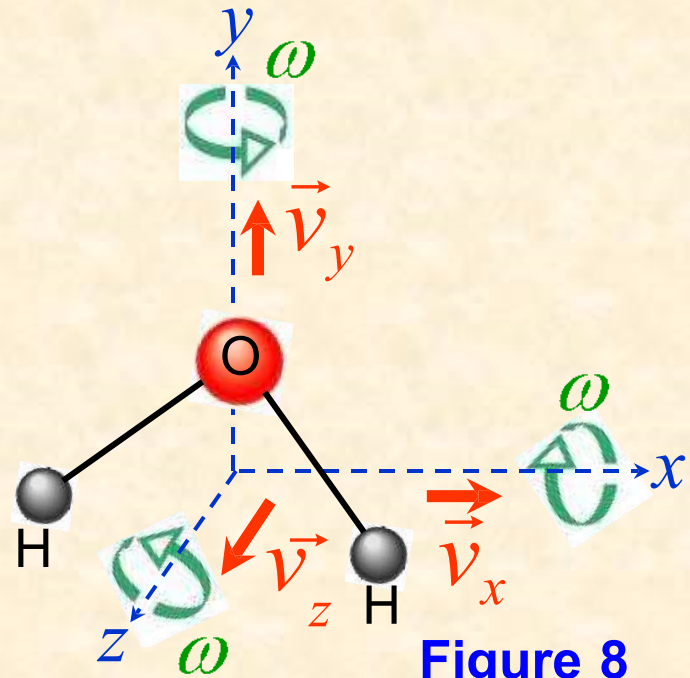


Figure 8

Table 1 shows the degrees of freedom for various molecules

Molecule	Example	Degrees of Freedom (f)			Average kinetic energy per molecule, $\langle K \rangle$
		Translational	Rotational	Total	
Monatomic	He	3	0	3	$\frac{3}{2}kT$
Diatomic	H ₂	3	2	5	$\frac{5}{2}kT$
Polyatomic	H ₂ O	3	3	6	$\frac{6}{2}kT = 3kT$

(At temperature of 300 K)

- Degrees of freedom depend on the **absolute temperature of the gases**.
 - For example : Diatomic gas (H_2)
 - Hydrogen gas have the vibrational kinetic energy (as shown in Figure 9) where contribute **2 degrees of freedom** which correspond to the **kinetic energy and the potential energy associated with vibrations along the bond** between the atoms.

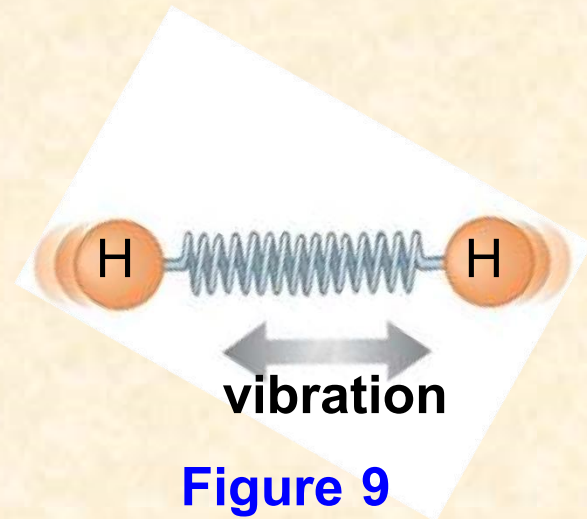


Figure 9

when the temperature,

At 250 K	→ $f = 3$
At 250 – 750 K	→ $f = 5$
At >750 K	→ $f = 7$

4 Internal energy of gas and relate the internal energy to the number of degree of freedom

- is defined as the **sum of total kinetic energy and total potential energy of the gas molecules**.
- But in ideal gas, the intermolecular forces are assumed to be negligible thus the **potential energy** of the molecules can be **neglected**.

Thus for N molecules,

$$U = N \langle K \rangle$$

$$U = \frac{f}{2} NkT$$

and

$$k = \frac{R}{N_A}$$

OR

$$U = \frac{f}{2} nRT$$

where U : internal energy
of the gas

Table 2 shows the properties for 1 mole of an ideal gas

	Monatomic	Diatomic	Polyatomic
Degrees of freedom, f	3	5	6
Average kinetic energy per molecule, $\langle K \rangle$	$\frac{3}{2}kT$	$\frac{5}{2}kT$	$\frac{6}{2}kT = 3kT$
Internal energy, U	$\frac{3}{2}RT$	$\frac{5}{2}RT$	$\frac{6}{2}RT = 3RT$